Turbulent drag reduction with streamwise travelling waves in the compressible regime

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Skin friction drag reduction with spanwise forcing

Streamwise-travelling waves of spanwise wall velocity

 $W(x,t) = Asin(\kappa_{x}x - \omega t)$

Oscillating wall: $\kappa_x = 0$ Steady wave: $\omega = 0$

Drag reduction(\mathcal{DR}) = 1 - $\frac{C_f}{C_{f,0}}$





 $Re_{ au} = 200, A^+ = 12, DR_{max} = 48\%$ Quadrio et al. JFM 2009

Towards real world applications

Dependence on:

• Reynolds number



• Complex shape and other drag sources



Quadrio et al. JFM 2022

• Mach number

≈ Yao & Hussain 2019

- Direct Numerical Simulations
- Channel flow
- STREAmS solver (Bernardini et al. CPC 2021)
- Constant flow rate (CFR) : $U_b = const$
- $Re_{\tau} = 400$
- Parameters:
 - Mach: $M_w^b = U_b/c_w = U_b/\sqrt{\gamma RT_w} = 0.3, 0.8, 1.5$ (as Yao & Hussain JFM 2019)
 - Control: $A^+ = 12$, 42 combinations of (ω^+, κ_y^+)

Simulation: control map



Results: Travelling waves at $\kappa_{\chi} = 0.005$



 $DR_{M=0.3} \approx 40 \%$ $DR_{M=1.5} \approx 52 \%$

Compressibility effect $(M_w^b \uparrow)$:

- Negative effect at large $|\omega|$
- Negative effect in the drag increase zone
- Large positive effect at small ω

- Travelling waves still work in compressible regime
- Compressibility has a positive effect on \mathcal{DR}

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- Compressibility has a positive effect on \mathcal{DR}

Is the increase of \mathcal{DR} actually due to compressibility?

The effect of the temperature

• $T_b \uparrow$ with control depending on (ω, κ_x)

$$T_b = \frac{1}{2h\rho_b U_b} \int_{-h}^{h} \langle \rho uT \rangle \ dy$$



The effect of the temperature

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Compressibility or inderect T_b \uparrow ?



The effect of the temperature

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Compressibility or inderect T_b \uparrow ?

• *T* profile not representative of external flows





$$\frac{\partial \rho e}{\partial t} + \frac{\partial \rho (e + p/\rho) u_j}{\partial x_j} = \frac{\partial \sigma_{ij} u_i}{\partial x_j} - \frac{\partial q_j}{\partial x_j} + f u_1$$

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Zero Bulk Cooling (ZBC)

$$\frac{\partial \rho e}{\partial t} + \frac{\partial \rho (e + p/\rho) u_j}{\partial x_j} = \frac{\partial \sigma_{ij} u_i}{\partial x_j} - \frac{\partial q_j}{\partial x_j} + f u_1 - \Phi$$

Zero Bulk Cooling (ZBC) Φ=0 Constained Bulk Cooling (CBC) $\Phi=f(\Theta=const)$

$$\frac{\partial \rho e}{\partial t} + \frac{\partial \rho (e + p/\rho) u_j}{\partial x_j} = \frac{\partial \sigma_{ij} u_i}{\partial x_j} - \frac{\partial q_j}{\partial x_j} + f u_1 - \Phi$$

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Diabatic Parameter (Θ)

$$\Theta = \frac{T_w - T_b}{T_r - T_b} \qquad T_r = \left(1 + \frac{\gamma - 1}{2}r(M_w^b)^2\right)T_b$$

 Θ = Fraction of the available kinetic energy transformed into thermal energy at the wall

We decouple compressibility from purely thermodynamical effects

Zero Bulk Cooling (ZBC) Θ varies

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T_b varies

Constained Bulk Cooling (CBC) ⊖ constant = 0.75 ↓ T_b

- constant across control
- varies across Mach

Drag reduction: ZBC vs CBC



 $DR_{M=0.3} \thickapprox 40 \% \quad DR_{M=1.5}^{ZBC} \thickapprox 52 \% \quad DR_{M=1.5}^{CBC} \thickapprox 43 \%$

Conclusions

- Influence of compressibility to the drag reduction performance of the streamwise-travelling waves of spanwise wall motion
- Two different comparison strategies:
 - Zero Bulk Cooling: ⊖=variable
 - Constrained Bulk Cooling: ⊖=const to decouple compressibility from purely thermodynamic effects

- Travelling waves still work in the compressible regime
- Compressibility effectiveness:
 - ZBC: large improvement
 - CBC: marginal improvement

Extension: Power budget

Drag reduction(\mathcal{DR})

$$\mathcal{DR}=1-\frac{C_f}{C_{f,0}}.$$

Input/control power (*P*_{in})

$$P_{in} = \frac{1}{P_0^*} \frac{1}{T_{ave} L_x L_z} \int_{t_i}^{t_f} \int_0^{L_x} \int_0^{L_z} W \, \tau_z \, dx \, dz \, dt$$

Pumping Power

$$P^* = \frac{U_b}{T_{ave}L_xL_z} \int_{t_i}^{t_f} \int_0^{L_x} \int_0^{L_z} \tau_x \, dx \, dz \, dt$$

Net energy saving rate (Pnet)

$$P_{net} = \mathcal{DR} - P_{in}.$$





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Results: Power budgets

Input/control power



 $P_{in} = \frac{1}{P_{0}^{*}} \frac{1}{T_{ave}L_{x}L_{z}} \int_{t_{i}}^{t_{f}} \int_{0}^{L_{x}} \int_{0}^{L_{z}} W \, \tau_{z} \, dx \, dz \, dt$





 $P_{not} = \mathcal{DR} - P_{in}$

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